

Corrections to Dashen's theorem from lattice QCD and quenched QED

Bálint C. Tóth

University of Wuppertal

Budapest–Marseille–Wuppertal-collaboration

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Outline

- 1 Introduction
- 2 QCD+QED
- 3 Definitions
- 4 Computational details
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Dashen's Theorem

- According to Dashen's theorem, in the $SU(3)$ limit [Dashen, 1969]

$$\Delta_{\text{QED}} M_K^2 = \Delta_{\text{QED}} M_\pi^2$$

where $\Delta_{\text{QED}} M_K^2 = (M_{K^-}^2 - M_{K^0}^2)_{\text{QED}}$ $\Delta_{\text{QED}} M_\pi^2 = (M_{\pi^-}^2 - M_{\pi^0}^2)_{\text{QED}}$

- Corrections are large away from $SU(3)$ limit
- Correction can be characterized by

$$\epsilon = \frac{\Delta_{\text{QED}} M_K^2 - \Delta_{\text{QED}} M_\pi^2}{\Delta M_\pi^2}$$

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QCD+QED

- Noncompact photon action

$$S_\gamma[A] = \frac{1}{4} \sum_{\mu, \nu, \mathbf{x}} (\partial_\mu A_{\mathbf{x}, \nu} - \partial_\nu A_{\mathbf{x}, \mu})^2$$

- Gauge fixing: Coulomb gauge $\underline{\partial}^\dagger \cdot \mathbf{A}_\mathbf{x} = 0$
- Zero mode subtraction: QED_L [Hayakawa & Uno, 2008]

$$\sum_{\mathbf{x}} A_{\mathbf{x}, t, \mu} = 0 \quad \text{for all } t \text{ and } \mu$$

- Coupling to quarks

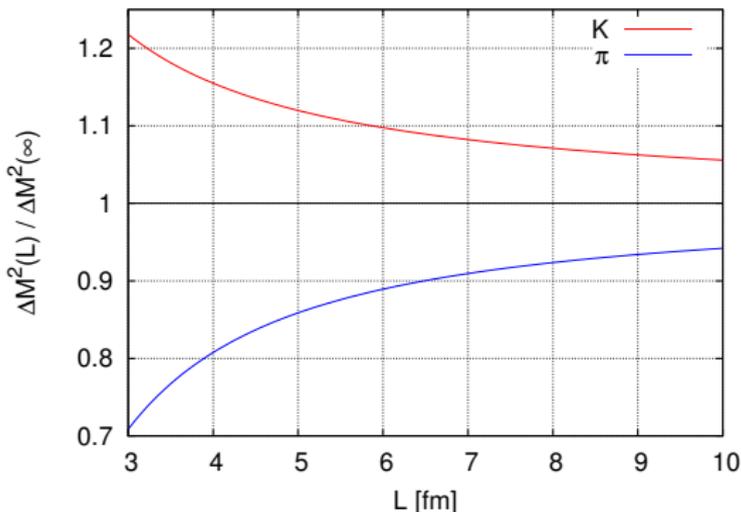
$$U_{\mathbf{x}, \mu} \longrightarrow U_{\mathbf{x}, \mu} \cdot e^{ieq A_{\mathbf{x}, \mu}}$$

FV correction in QED_L

$$m(T, L) \sim m \cdot \left(1 - \frac{q^2 \alpha \kappa}{2 m L} \left[1 + \frac{2}{m L} \right] + \mathcal{O}\left(\frac{\alpha}{L^3}\right) \right)$$

with $\kappa \approx 2.837297$

[Borsanyi *et.al.*, 2015]



Quenched QED

- QCD + QED

$$\langle \mathcal{O} \rangle_{U,A} = \frac{1}{Z} \int dU \int dA \mathcal{O}(U, A) e^{-S_g[U]} e^{-S_\gamma[A]} \det M(U, A)$$

- Quenched QED

$$\begin{aligned} \langle \mathcal{O} \rangle_{\substack{U, \text{dynamical} \\ A, \text{quenched}}} &= \frac{1}{Z} \int dU \int dA \underbrace{\mathcal{O}(U, A) e^{-S_\gamma[A]} e^{-S_g[U]} \det M(U)}_{\langle \mathcal{O}(U, A) \rangle_{A,q}} \\ &= \left\langle \left\langle \mathcal{O}(U, A) \right\rangle_{A,q} \right\rangle_U \end{aligned}$$

- In practice

- Generate $SU(3)$ configurations
- For each $SU(3)$ configuration, generate $U(1)$ configurations with S_γ
- Measure $\mathcal{O}(U, A)$

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Physical point

bare parameters	→	measurable observables
$m_l = \frac{m_d + m_u}{2}$		$M_{\pi^+}^2$
$\delta m = m_d - m_u$		$\Delta M_K^2 = M_{K^0}^2 - M_{K^+}^2$
m_s		$M_{K^X}^2 = M_{K^0}^2 + M_{K^+}^2 - M_{\pi^+}^2$
α		α
a		M_Ω

- We are at the physical point, if

$$\frac{M_{\pi^+}^2}{M_\Omega^2} = \left(\frac{M_{\pi^+}^2}{M_\Omega^2} \right)_{\text{phys}}, \quad \frac{\Delta M_K^2}{M_\Omega^2} = \left(\frac{\Delta M_K^2}{M_\Omega^2} \right)_{\text{phys}}, \quad \frac{M_{K^X}^2}{M_\Omega^2} = \left(\frac{M_{K^X}^2}{M_\Omega^2} \right)_{\text{phys}}, \quad \alpha = \alpha_{\text{phys}}$$

- In vicinity of physical point

$$\begin{aligned} \mathcal{O}(\{m_l, \delta m, m_s, \alpha\}) &= \mathcal{O}_{\text{phys}} + B \left(\frac{M_{\pi^+}^2}{M_\Omega^2} - \left(\frac{M_{\pi^+}^2}{M_\Omega^2} \right)_{\text{phys}} \right) + C \left(\frac{\Delta M_K^2}{M_\Omega^2} - \left(\frac{\Delta M_K^2}{M_\Omega^2} \right)_{\text{phys}} \right) + \\ &+ D \left(\frac{M_{K^X}^2}{M_\Omega^2} - \left(\frac{M_{K^X}^2}{M_\Omega^2} \right)_{\text{phys}} \right) + E (\alpha - \alpha_{\text{phys}}) \end{aligned}$$

m_u/m_d

- Introduce

M_{uu} : meson with quark content u, \bar{u} , w/o disconnected

M_{dd} : meson with quark content d, \bar{d} , w/o disconnected

- Connection to quark masses

$$\begin{aligned}
 M_{uu}^2 &\propto m_u \\
 M_{dd}^2 &\propto m_d \\
 \Delta M^2 = M_{dd}^2 - M_{uu}^2 &\propto m_d - m_u \\
 M_{\pi^0}^2 &\propto \frac{m_d + m_u}{2}
 \end{aligned}$$

- Define

$$\frac{m_u}{m_d} = \frac{2M_{\pi^0}^2 - M_{dd}^2 + M_{uu}^2}{2M_{\pi^0}^2 + M_{dd}^2 - M_{uu}^2} = \frac{2M_{\pi^0}^2 - \Delta M^2}{2M_{\pi^0}^2 + \Delta M^2}$$

Strategy to obtain m_u/m_d

- Measure $\Delta M^2 = M_{dd}^2 - M_{uu}^2$ in continuum, at physical point

$$\left(\frac{\Delta M^2}{M_\Omega^2}\right)(\{m_l, \delta m, m_s, \alpha\}) = \left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{phys}} + B \left(\frac{M_{\pi^+}^2}{M_\Omega^2} - \left(\frac{M_{\pi^+}^2}{M_\Omega^2}\right)_{\text{phys}}\right) + C \left(\frac{\Delta M_K^2}{M_\Omega^2} - \left(\frac{\Delta M_K^2}{M_\Omega^2}\right)_{\text{phys}}\right) + D \left(\frac{M_{K^*}^2}{M_\Omega^2} - \left(\frac{M_{K^*}^2}{M_\Omega^2}\right)_{\text{phys}}\right) + E(\alpha - \alpha_{\text{phys}})$$

- Then

$$\frac{m_u}{m_d} = \frac{2 \left(\frac{M_{\pi^0}^2}{M_\Omega^2}\right)_{\text{phys}} - \left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{phys}}}{2 \left(\frac{M_{\pi^0}^2}{M_\Omega^2}\right)_{\text{phys}} + \left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{phys}}}$$

Correction to Dashen's Theorem

- Correction is characterised by

$$\epsilon = \frac{\Delta_{\text{QED}} M_K^2 - \Delta_{\text{QED}} M_\pi^2}{\Delta M_\pi^2}$$

where $\Delta_{\text{QED}} M_K^2 = (M_{K^-}^2 - M_{K^0}^2)_{\text{QED}}$ $\Delta_{\text{QED}} M_\pi^2 = (M_{\pi^-}^2 - M_{\pi^0}^2)_{\text{QED}}$

- $\Delta_{\text{QED}} M_\pi^2$ is difficult to measure

ΔM_π^2 is dominated by $\Delta_{\text{QED}} M_\pi^2$

[FLAG, 2017]

- Measure instead

$$\epsilon' = \frac{\Delta_{\text{QED}} M_K^2 - \Delta M_\pi^2}{\Delta M_\pi^2}$$

- We need to obtain $\Delta_{\text{QED}} M_K^2$ in continuum, at physical point

Strategy to obtain $\Delta_{\text{QED}} M_K^2$

- $\frac{\Delta M_K^2}{M_\Omega^2} = \left(\frac{\Delta_{\text{QED}} M_K^2}{M_\Omega^2} \right)$ if

$$\left(\frac{\Delta M^2}{M_\Omega^2} \right) = 0, \quad \frac{M_{\pi^+}^2}{M_\Omega^2} = \left(\frac{M_{\pi^+}^2}{M_\Omega^2} \right)_{\text{phys}}, \quad \frac{M_{K^X}^2}{M_\Omega^2} = \left(\frac{M_{K^X}^2}{M_\Omega^2} \right)_{\text{phys}}, \quad \alpha = \alpha_{\text{phys}}$$

- Measure $\left(\frac{\Delta M^2}{M_\Omega^2} \right)_{\text{phys}}$ and C

$$\begin{aligned} \left(\frac{\Delta M^2}{M_\Omega^2} \right) &= \left(\frac{\Delta M^2}{M_\Omega^2} \right)_{\text{phys}} + B \left(\frac{M_{\pi^+}^2}{M_\Omega^2} - \left(\frac{M_{\pi^+}^2}{M_\Omega^2} \right)_{\text{phys}} \right) + \\ &+ C \left(\frac{\Delta M_K^2}{M_\Omega^2} - \left(\frac{\Delta M_K^2}{M_\Omega^2} \right)_{\text{phys}} \right) + D \left(\frac{M_{K^X}^2}{M_\Omega^2} - \left(\frac{M_{K^X}^2}{M_\Omega^2} \right)_{\text{phys}} \right) + E (\alpha - \alpha_{\text{phys}}) \end{aligned}$$

- Obtain $\frac{\Delta_{\text{QED}} M_K^2}{M_\Omega^2}$ from

$$0 = \left(\frac{\Delta M^2}{M_\Omega^2} \right)_{\text{phys}} + C \left(\frac{\Delta_{\text{QED}} M_K^2}{M_\Omega^2} - \left(\frac{\Delta M_K^2}{M_\Omega^2} \right)_{\text{phys}} \right)$$

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Ensembles

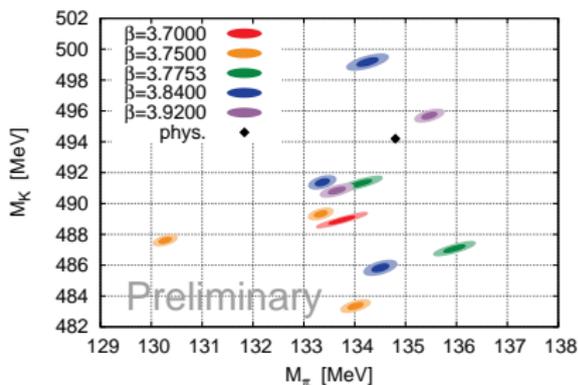
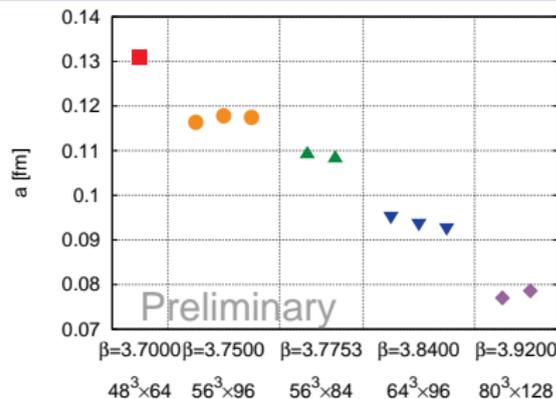
- Subset of [1612.02364,1711.04980]
- Tree-level Symanzyk gauge action
- $N_f = 2 + 1 + 1$ staggered fermions
- stout smearing 4 steps, $\varrho = 0.125$
- m_l and m_s are chosen to give approx.

$$\overline{M}_\pi = 134.8(3) \text{ MeV}$$

$$\overline{M}_K = 494.2(3) \text{ MeV}$$

[FLAG, 2017]

- m_c is fixed via $\frac{m_c}{m_s} = 11.85$
- α is set to Thomson limit
- a is fixed via $M_\Omega = 1672.45 \text{ MeV}$



Fit strategy

- Measurements at 3 parameter sets:

$$\text{ISO: } \delta m = 0, \quad \alpha = 0 \quad \longrightarrow \quad \Delta M^2 = 0, \quad \Delta M_K^2 = 0$$

$$\text{QED: } \delta m = 0, \quad \alpha = \alpha_{\text{phys}}$$

$$\text{SIB: } \delta m \neq 0, \quad \alpha = 0$$

- Fully correlated fit for $\left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{phys}}$, A_γ , A_s , B , C , D

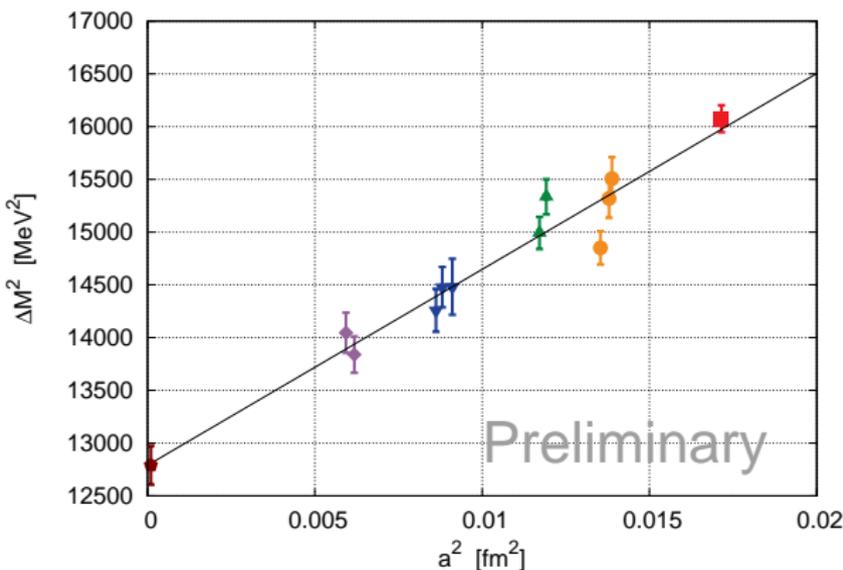
$$\begin{aligned} \left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{QED}} &= \left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{phys}} + A_\gamma \cdot (aM_\Omega)^2 + B \left(\frac{M_{\pi^+}^2}{M_\Omega^2}(\text{QED}) - \left(\frac{M_{\pi^+}^2}{M_\Omega^2}\right)_{\text{phys}} \right) + \\ &+ C \left(\frac{\Delta M_K^2}{M_\Omega^2}(\text{QED}) - \left(\frac{\Delta M_K^2}{M_\Omega^2}\right)_{\text{phys}} \right) + D \left(\frac{M_{K^*}^2}{M_\Omega^2}(\text{QED}) - \left(\frac{M_{K^*}^2}{M_\Omega^2}\right)_{\text{phys}} \right) \end{aligned}$$

$$\begin{aligned} \left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{SIB}} - \left(\frac{\Delta M^2}{M_\Omega^2}\right)_{\text{ISO}} &= A_s \cdot (aM_\Omega)^2 + \\ &+ B \left(\frac{M_{\pi^+}^2}{M_\Omega^2}(\text{SIB}) - \frac{M_{\pi^+}^2}{M_\Omega^2}(\text{ISO}) \right) + C \left(\frac{\Delta M_K^2}{M_\Omega^2}(\text{SIB}) - \frac{\Delta M_K^2}{M_\Omega^2}(\text{ISO}) \right) \end{aligned}$$

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Continuum limit

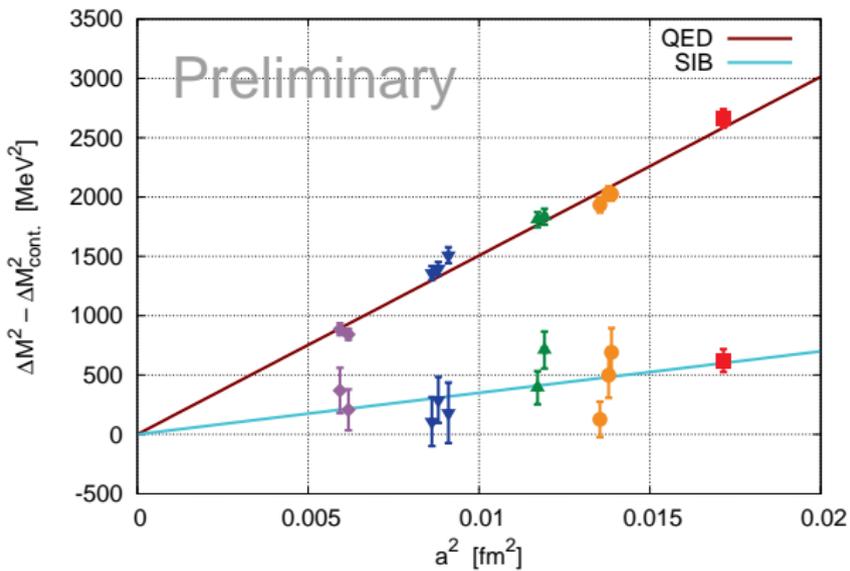


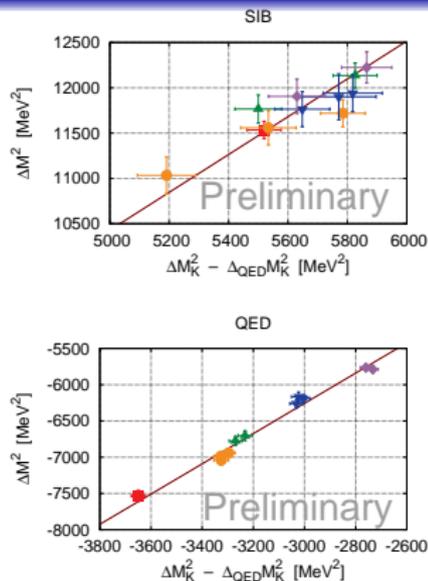
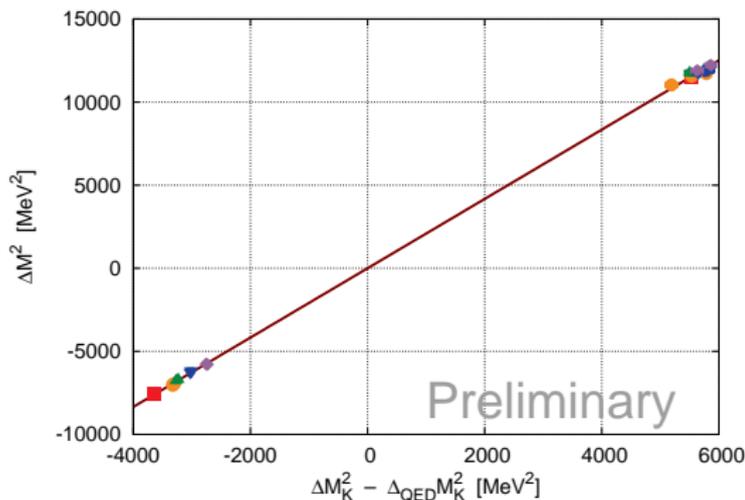
$$\Delta M^2 = 1.279(18)(17) \times 10^4 \text{ MeV}^2$$

$$\frac{m_u}{m_d} = \frac{2M_{\pi_0}^2 - \Delta M^2}{2M_{\pi_0}^2 + \Delta M^2} = 0.4804(55)(52)$$

A_γ

vs.

 A_S 

$\Delta_{\text{QED}} M_K^2$ 

$$C = 2.085(29)(12)$$

$$\Delta M^2 = \Delta M_{\text{phys}}^2 + C (\Delta M_K^2 - \Delta M_{K,\text{phys}}^2) + \dots$$

$$\Delta_{\text{QED}} M_K^2 = \Delta M_K^2 - \frac{1}{C} \cdot \Delta M^2 = 2228(35)(47) \text{ MeV}^2$$

$$\epsilon' = \frac{\Delta_{\text{QED}} M_K^2 - \Delta M_\pi^2}{\Delta M_\pi^2} = 0.767(28)(38)$$

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Summary

- Computation of m_u/m_d and $\Delta_{\text{QED}} M_K^2$ using $N_f = 2 + 1 + 1$ staggered fermions
- QED is still quenched
- Preliminary results, compatible with previous calculations